Differential Kinematics

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Differential Kinematics

Up to this point, we have only considered the relationship of the joint angles to the Cartesian location of the end effector:

f(q) = x

But what about the first derivative?



• This would tell us the velocity of the end effector as a function of joint angle velocities.

Motivating Example

Consider a one-link arm

- As the arm rotates, the end effector sweeps out an arc
- Let's assume that we are only interested in the *x* coordinate...



Differential kinematics: $\frac{dx}{dq} = -l\sin(q)$

$$\delta x = -l\sin(q)\delta q$$
$$\delta q = -\frac{1}{l\sin(q)}\delta x$$



Motivating Example

Suppose you want to move the end effector above a specified point, X_a

Answer #1:
$$q_g = \cos^{-1}\left(\frac{x_g}{l}\right)$$

Answer #2: 1. $i = 0, q_0$ = arbitrary

2.
$$x_i = l\cos(q_i)$$

3. $\delta x = \alpha (x_g - x_i)$
4. $\delta q = \frac{1}{-l\sin(q_i)} \delta x$

5. $q_{i+1} = q_i + \delta q$ 6. i + + goto 2.





This controller moves the link asymptotically toward the goal position.

Intro to the Jacobian

$$x = \begin{bmatrix} l_1 \cos(q_1) + l_2 \cos(q_1 + q_2) \\ l_1 \sin(q_1) + l_2 \sin(q_1 + q_2) \end{bmatrix} \longleftarrow \qquad \text{Forward kinematics of the two-link manipulator}$$

$$Velocity \text{ Jacobian}$$

$$\downarrow$$

$$\frac{dx}{dq} = \begin{pmatrix} -l_1 \sin(q_1) - l_2 \sin(q_1 + q_2) \\ l_1 \cos(q_1) + l_2 \cos(q_1 + q_2) \end{bmatrix} - l_2 \sin(q_1 + q_2) \\ -l_2 \cos(q_1 + q_2) \end{pmatrix}$$

$$= J[q]$$

Intro to the Jacobian



Jacobian

The Jacobian relates joint velocities with end effector *twist*:



It turns out that you can "easily" compute the Jacobian for arbitrary manipulator structures

• This makes differential kinematics a much easier sub-problem than kinematics in general.

What is Twist?

End effector twist:

- Twist is a concatenation of linear velocity and angular velocity:
- As we will show in a minute, linear and angular velocity have different units



What is Twist?



Angular velocity is a vector that:

- points in the direction of the axis of rotation
- has magnitude equal to the velocity of rotation

What is Angular Velocity?

Angular velocity is a vector that:

- points in the direction of the axis of rotation

- has magnitude equal to the velocity of rotation

Symbol for angular velocity: $\,\omega\,$

Relation between angular velocity and linear velocity: $v=\omega\times r$



We will often write it this way: $\,\dot{q}=\omega imes q\,$

Angular Velocity Derivation

$${}^{b}q \stackrel{a}{=} R_{a}{}^{a}q$$

 ${}^{b}\dot{q} \stackrel{b}{=} \dot{R}_{a}{}^{a}q \leftarrow$

 ${}^{b}\dot{q} = {}^{b}\dot{R}_{a}{}^{b}R_{a}{}^{T}{}^{b}q$

Just differentiate all elements of the rotation matrix w.r.t. time.



This is the matrix representation of angular velocity

This FO differential equation encodes how the particle rotates

Twist: Time out for skew symmetry!

$$S = -S^{T} \qquad \longrightarrow \qquad \text{Def'n of skew symmetry}$$

$$S = \begin{bmatrix} 0 & a & b \\ -a & 0 & c \\ -b & -c & 0 \end{bmatrix} \qquad \longleftarrow \qquad \text{Skew symmetric matrices}$$
always look like this
If you interpret the skew symmetric
$$S(x) = \begin{bmatrix} 0 & -x_{z} & x_{y} \\ x_{z} & 0 & -x_{x} \\ -x_{y} & x_{x} & 0 \end{bmatrix}$$

Then this is another way of writing the cross product:

 $S(x)p = x \times p$

Angular Velocity Derivation

Skew symmetry of $S^{(b_{\omega})}$:

$$I \stackrel{b}{=} \stackrel{b}{R}_{a} \stackrel{b}{R}_{a} \stackrel{T}{=} \stackrel{c}{=} \stackrel{b}{R}_{a} \stackrel{b}{R}_{a} \stackrel{T}{=} \stackrel{c}{=} \stackrel{b}{R}_{a} \stackrel{b}{R}_{a} \stackrel{T}{=} \stackrel{c}{=} \stackrel{b}{\omega} \stackrel{b}{q} \stackrel{c}{=} \stackrel{c}{=} \stackrel{c}{\omega} \stackrel{b}{\omega} \stackrel{c}{=} \stackrel{c}{=} \stackrel{c}{\omega} \stackrel{b}{\omega} \stackrel{c}{=} \stackrel{c}{=} \stackrel{c}{\omega} \stackrel{b}{\omega} \stackrel{c}{=} \stackrel{c}{=} \stackrel{c}{\omega} \stackrel{b}{=} \stackrel{c}{=} \stackrel{c}{\omega} \stackrel{b}{=} \stackrel{c}{=} \stackrel{c}{\omega} \stackrel{c}{=} \stackrel{c}{=} \stackrel{c}{\omega} \stackrel{c}{=} \stackrel{c}{=} \stackrel{c}{\omega} \stackrel{c}{=} \stackrel{c}{=} \stackrel{c}{\omega} \stackrel{c}{=} \stackrel{c}{=} \stackrel{c}{=} \stackrel{c}{\omega} \stackrel{c}{=} \stackrel{$$

You probably already know this formula

Twist

Twist concatenates linear and angular velocity:



Jacobian

Breakdown of the Jacobian: $v = J_v q$ $\omega = J_{\omega}\dot{q}$ $\xi = \begin{bmatrix} J_{\nu} \\ J_{\omega} \end{bmatrix} \dot{q}$ Relation to the derivative: $J_v = \frac{\partial x}{\partial q}$ but $J_\omega \neq \frac{\partial r_{\phi\theta\psi}}{\partial q}$

That's not an angular velocity

Calculating the Jacobian

Approach:

- Calculate the Jacobian one column at a time
- Each column describes the motion at the end effector due to the motion of *that joint only.*
- For each joint, *i*, pretend all the other joints are frozen, and calculate the motion at the end effector caused by *i*.



Calculating the Jacobian: Velocity



Velocity at end effector due to rotation at joint *i-1*

Calculating the Jacobian: Velocity



 $J_{v_i} \stackrel{b}{=} z_{i-1}$

Extension/contraction along $i^{-1}z$

Calculating the Jacobian: Velocity

Rotational DOF

• Rotates about $i-1_Z$

$$J_{\omega_i} \stackrel{b}{=} z_{i-1,i}$$



Prismatic DOF

• Translates along $i-1_Z$

 $J_{\omega_i} = 0$

Extension/contraction along $i^{-1}z$

Calculating the Jacobian: putting it together $J_{v} = \begin{bmatrix} J_{v_{1}} & \cdots & J_{v_{n}} \end{bmatrix}$

Where

- rotational $J_{v_i} \stackrel{b}{=} z_{i-1} \times ({}^b p_{eff} {}^b p_{i-1})$ prismatic $J_{v_i} \stackrel{b}{=} z_{i-1}$

$$J_{\omega} = \begin{bmatrix} J_{\omega_1} & \cdots & J_{\omega_n} \end{bmatrix}$$

Where

- rotational $J_{\omega_i} \stackrel{p}{=} z_{i-1}$
- prismatic $J_{\omega} = 0$

$$J = \begin{bmatrix} J_{v_1} & \cdots & J_{v_n} \\ J_{\omega_1} & \cdots & J_{\omega_n} \end{bmatrix}$$



From before:

$${}^{0}T_{1} = \begin{pmatrix} c_{q_{1}} & -s_{q_{1}} & 0 & l_{1}c_{q_{1}} \\ s_{q_{1}} & c_{q_{1}} & 0 & l_{1}s_{q_{1}} \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \qquad {}^{1}T_{2} = \begin{pmatrix} c_{q_{2}} & -s_{q_{2}} & 0 & l_{2}c_{q_{2}} \\ s_{q_{2}} & c_{q_{2}} & 0 & l_{2}s_{q_{2}} \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

$${}^{2}T_{3} = \begin{pmatrix} c_{q_{3}} & -s_{q_{3}} & 0 & l_{3}c_{q_{3}} \\ s_{q_{3}} & c_{q_{3}} & 0 & l_{3}s_{q_{3}} \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

$$J_{\omega} = \begin{bmatrix} 0 \hat{z}_0 & 0 \hat{z}_1 & 0 \hat{z}_2 \end{bmatrix} = \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 1 & 1 & 1 \end{bmatrix}$$



$$\begin{split} J_{v_{1}} &= {}^{0} \hat{z}_{0} \dot{\iota} \left({}^{0} o_{3} - {}^{0} o_{0} \right) = \begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix} \times \begin{pmatrix} \begin{bmatrix} l_{1} c_{1} + l_{2} c_{12} + l_{3} c_{123} \\ l_{1} s_{1} + l_{2} s_{12} + l_{3} s_{123} \\ 0 \end{bmatrix} - \begin{bmatrix} 0 \\ 0 \\ 0 \\ 0 \end{bmatrix} \end{pmatrix} = \begin{bmatrix} -l_{1} s_{1} - l_{2} s_{12} - l_{3} s_{123} \\ l_{1} c_{1} + l_{2} c_{12} + l_{3} c_{123} \\ 0 \end{bmatrix} \\ J_{v_{2}} &= {}^{0} \hat{z}_{1} \dot{\iota} \left({}^{0} o_{3} - {}^{0} o_{1} \right) = \begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix} \times \begin{pmatrix} \begin{bmatrix} l_{1} c_{1} + l_{2} c_{12} + l_{3} c_{123} \\ l_{1} s_{1} + l_{2} s_{12} + l_{3} s_{123} \\ 0 \end{bmatrix} - \begin{bmatrix} l_{1} c_{1} \\ l_{1} s_{1} \\ 0 \end{bmatrix} \end{pmatrix} = \begin{bmatrix} -l_{2} s_{12} - l_{3} s_{123} \\ l_{2} c_{12} + l_{3} c_{123} \\ l_{2} c_{12} + l_{3} c_{123} \\ 0 \end{bmatrix} \\ J_{v_{3}} &= {}^{0} \hat{z}_{2} \dot{\iota} \left({}^{0} o_{3} - {}^{0} o_{2} \right) = \begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix} \times \begin{pmatrix} \begin{bmatrix} l_{1} c_{1} + l_{2} c_{12} + l_{3} c_{123} \\ l_{1} s_{1} + l_{2} s_{12} + l_{3} s_{123} \\ 0 \end{bmatrix} - \begin{bmatrix} l_{1} c_{1} + l_{2} c_{12} \\ l_{1} s_{1} + l_{2} s_{12} \\ 0 \end{bmatrix} = \begin{bmatrix} -l_{3} s_{123} \\ l_{3} c_{123} \\ l_{3} c_{123} \\ 0 \end{bmatrix} \end{split}$$

³y

 ^{2}v

х

$$J_{v} = \begin{bmatrix} -l_{1}s_{1} - l_{2}s_{12} - l_{3}s_{123} & -l_{2}s_{12} - l_{3}s_{123} & -l_{3}s_{123} \\ l_{1}c_{1} + l_{2}c_{12} + l_{3}c_{123} & l_{2}c_{12} + l_{3}c_{123} & l_{3}c_{123} \\ 0 & 0 & 0 \end{bmatrix}$$

$$J = \begin{bmatrix} J_{\nu} \\ J_{\omega} \end{bmatrix} = \begin{bmatrix} -l_{1}s_{1} - l_{2}s_{12} - l_{3}s_{123} & -l_{2}s_{12} - l_{3}s_{123} & -l_{3}s_{123} \\ l_{1}c_{1} + l_{2}c_{12} + l_{3}c_{123} & l_{2}c_{12} + l_{3}c_{123} & l_{3}c_{123} \\ 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \\ 1 & 1 & 1 \end{bmatrix}$$

Think-pair-share



Calculate the end effector Jacobian with respect to the base frame







$$J_{v_{1}} = \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} \times \begin{pmatrix} -c_{1}(l_{2}c_{2}+l_{3}c_{23}) \\ -s_{1}(l_{2}c_{2}+l_{3}c_{23}) \\ l_{2}s_{2}+l_{3}s_{23}+l_{1} \end{pmatrix} = \begin{pmatrix} s_{1}(l_{2}c_{2}+l_{3}c_{23}) \\ -c_{1}(l_{2}c_{2}+l_{3}c_{23}) \\ 0 \end{pmatrix} \qquad J_{\omega_{1}} = \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}$$

$$J_{\omega_{1}} = \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}$$

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$$J_{\omega_{2}} = \begin{pmatrix} -s_{1} \\ c_{1} \\ 0 \end{pmatrix}$$

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$$J = \begin{pmatrix} s_1(l_2c_2+l_3c_{23}) & c_1(l_2c_2+l_3c_{23}) & l_3c_1s_{23} \\ -c_1(l_2c_2+l_3c_{23}) & s_1(l_2c_2+l_3c_{23}) & l_3c_1s_{23} \\ 0 & l_2c_2+l_3c_{23} & l_3c_{23} \\ 0 & -s_1 & -s_1 \\ 0 & c_1 & c_1 \\ 1 & 0 & 0 \end{pmatrix}$$

 $y_1 \qquad y_2 \qquad x_2 \qquad y_3 \qquad y_1 \qquad z_2 \qquad y_3 \qquad y_1 \qquad z_2 \qquad y_3 \qquad y_3 \qquad y_1 \qquad z_2 \qquad y_3 \qquad y_3 \qquad y_3 \qquad y_4 \qquad y_5 \qquad y_5$

Expressing the Jacobian in Different Reference Frames

In the preceeding, the Jacobian has been expressed in the base frame

• It can be expressed in other reference frames using rotation matrices

Velocity is transformed from one reference frame to another using:

$${}^{k} p = {}^{k} R_{b} {}^{b} p$$
$${}^{k} \dot{p} = {}^{k} R_{b} {}^{b} \dot{p}$$

Therefore, the velocity Jacobian can be transformed using: ${}^{k}J_{\nu} = {}^{k}R_{b}{}^{b}J_{\nu}$



Expressing the Jacobian in Different Reference Frames x^{3x}

First, let's express angular velocity in a different reference frame:



 ${}^{k}\omega = {}^{k}R_{b}{}^{b}\omega$ \leftarrow Angular velocity can also be rotated by a rotation matrix

Therefore, the angular velocity Jacobian can be transformed using:

$${}^{k}J_{\omega} = {}^{k}R_{b}{}^{b}J_{\omega}$$

Expressing the Jacobian in Different Reference Frames x^{3x}

Therefore, the full Jacobian is rotated:

$${}^{k}J = \begin{pmatrix} {}^{k}R_{b} & 0 \\ 0 & {}^{k}R_{b} \end{pmatrix} {}^{b}J$$



Different Jacobian Reference Frames: Example

Express the Jacobian for the three-link arm in the y - reference frame of the end effector:

$${}^{0}R_{3} = \begin{pmatrix} c_{123} & -s_{123} & 0 \\ s_{123} & c_{123} & 0 \\ 0 & 0 & 1 \end{pmatrix}$$
$$J = \begin{bmatrix} -l_{1}s_{1} - l_{2}s_{12} - l_{3}s_{123} & -l_{2}s_{12} - l_{3}s_{123} & -l_{3}s_{123} \\ l_{1}c_{1} + l_{2}c_{12} + l_{3}c_{123} & l_{2}c_{12} + l_{3}c_{123} & l_{3}c_{123} \\ 0 & 0 & 0 \\ 0 & 0 & 0 \\ 1 & 1 & 1 & 1 \end{bmatrix}$$



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Think-pair-share

